

# Can Nuclear Magnetic Resonance be Used as a Plasma Diagnostic ?

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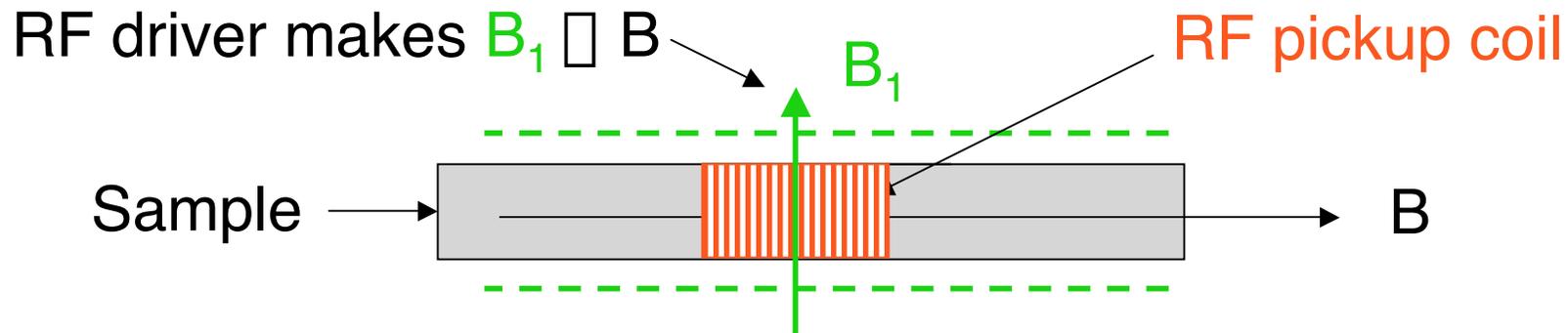
PS&T Seminar  
5/9/03

see also: Rev. Sci. Inst. 74, 11460 (2003)  
from High Temperature Plasma Diagnostics Conf. 2002

# Outline

- Conventional NMR in liquids and solids
- Plasma NMR vs. Conventional NMR
- Possible plasma NMR measurements
- Estimates of signal/noise ratios
- Conclusions and possible directions

# Conventional NMR in Liquids and Solids



- Insert sample into very uniform magnetic field  $B$  ( $\approx 1-20$  T), and apply transverse oscillating  $B_1$  pulse at NMR frequency to rotate nuclear magnetization  $\perp B$  (typically over  $\approx 10 \mu\text{sec}$ )
- After  $B_1$  field ends, measure precession of nuclear spins using RF pickup coil (“Free Induction Decay” phase)
- Frequency and amplitude of NMR spectrum in pickup coil determine type and number of nuclei in sample

# NMR Nuclei and Frequencies

- Nuclei used for NMR are of interest to fusion
- Frequencies are (coincidentally) in ICRH range

Particle	Spin	$f_{\text{NMR}}$ (MHz)*	$f_{\text{NMR}}/f_{\text{ci}}$ **	Sensitivity***
Hydrogen	1/2	100	2.8	1
Deuterium	1	15.35	0.86	0.009
Tritium	1/2	106.7	9.0	1.2
He-3	1/2	76.2	3.2	0.44
Li-6	1	14.7	0.82	0.009
Li-7	3/2	38.9	2.5	0.29
Beryllium-9	3/2	14.1	0.89	0.014
Boron-10	3	10.7	0.6	0.02
Boron-11	3/2	32.1	2	0.17
Carbon-13	1/2	25.1	1.5	0.016

\* at  $B=2.3488$  T, from Abraham et al, Introduction to NMR Spectroscopy, 1988

\*\* for fully ionized nuclei, equal to  $f_{\text{NMR}}/f_{\text{ci}} = (1/2)(\mu/\mu_N)(\mu_i/\mu_p)(1/Z)$

\*\*\* relative to protons at constant B for equal number of nuclei

# Examples of NMR Machines and Results

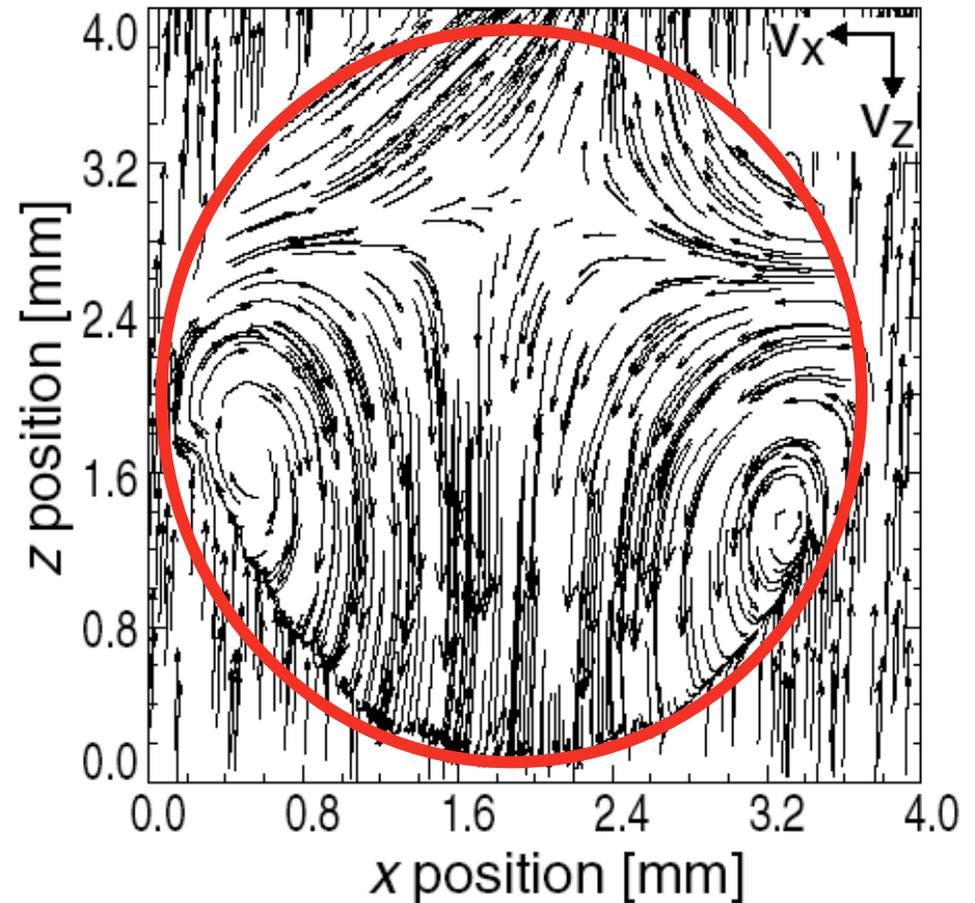
21 Tesla NMR machine



**A 900-MHz nuclear magnetic resonance device.**

Physics Today, Mar. 2002

velocity of falling water drop



Han et al, Phys. Rev. Lett. 87, 144501 (2001)

# Conventional NMR vs. Plasma NMR

- In general NMR signal  $\propto$  density of NMR species in sample
  - $\Rightarrow$  plasma signal  $\approx 10^{-9}$  times liquid signal if they have the same fractional spin polarization
- Spin polarization fraction typically  $F_{\text{spin}} \approx \mu B / kT_{\text{room}} \approx 10^{-5}$ 
  - $\Rightarrow$  in equilibrium expect  $F_{\text{spin}} \approx 10^{-10}$  at  $T = 2.5 \text{ keV}$
- However, nearly 100% polarized  $\text{H}_2$  or  $^3\text{He}$  can be made in large quantities for use as plasma “fuel” for NMR
  - $\Rightarrow$  fully polarized plasma would increase NMR signal by  $\approx 10^5$  with respect to room temperature sample

# Plasma NMR Signal with Polarized Fuel

- $V \approx 10^{-3}$  volt NMR signal from hydrogen in 1 cm<sup>3</sup> of water at room temperature with, B=1 T using a 10 turn coil with A=1 cm<sup>2</sup> at f=42 MHz (Bloch, Phys. Rev. p. 460, 1946)
- With respect to water, plasma density is smaller by  $\approx 10^{-9}$ , but with fully polarized H plasma, polarized spin density fraction would be larger by  $\approx 10^5$
- Thus if plasma area was A=1 m<sup>2</sup> instead of 1 cm<sup>2</sup>, NMR signal from plasma should be comparable to NMR signal from 1 cm<sup>3</sup> of water (measured in 1940's)

expect  $V_{\text{plasma}} \approx 10^{-3}$  volts  $\Rightarrow$  easily measurable ! ?

# Possible Sources of Polarized Plasma Fuel

- Highly polarized  $^3\text{He}$  gas ( $> 10^{21}$  atoms) produced by spin exchange with laser optically pumped alkali-metal vapor (e.g. Chupp et al, Phys. Rev. C, 2244 '87)
- Bottles of this gas can be transported and kept polarized for many hours if a weak magnetic field is present
- Similarly polarized hydrogen has also been created using cryogenic methods (Cline et al, PRL 47, 1195 '81)
- In principle, highly polarized beams or pellets can also be created and used for plasma fueling

# Spin Depolarization in Plasmas

Various processes calculated by Kulsrud et al in papers on spin-polarized fusion (e.g. Nucl. Fusion 186, p. 1443)

- Collisional depolarization negligible for fusion plasmas
- Effect of ion motion in static tokamak fields negligible
- Effect of magnetic field fluctuations near NMR frequency can cause rapid depolarization (similar to  $B_1$  field)  
Kulsrud et al estimated  $\tau_{\text{decay}} \approx .075$  sec for decay time of spin polarized tritium for assumed plasma  $\Delta B_{\parallel} \approx 1$  Gauss at  $B=5T$
- Effect of nuclei recycling from wall can also be significant  
typically polarization lost in  $\leq 1$  sec when atom resides on wall surface  
H.S. Greenside et al, J. Vac. Sci. Technol. A 2, 619 (1984)

# Possible NMR Experiments on Tokamaks

## 1) Plasma NMR

- Measure species mix inside a fully ionized plasma  
(but needs polarized fuel for each species !)
- Measure internal  $\beta$  from decay rate of polarization

## 2) Wall surface NMR measurement (without plasma)

- Tritium inventory inside vessel (typically  $\approx 1$  gram in JET)
- Low-Z surface coatings used for conditioning (Li, Be, B)

# How Could This be Done ?

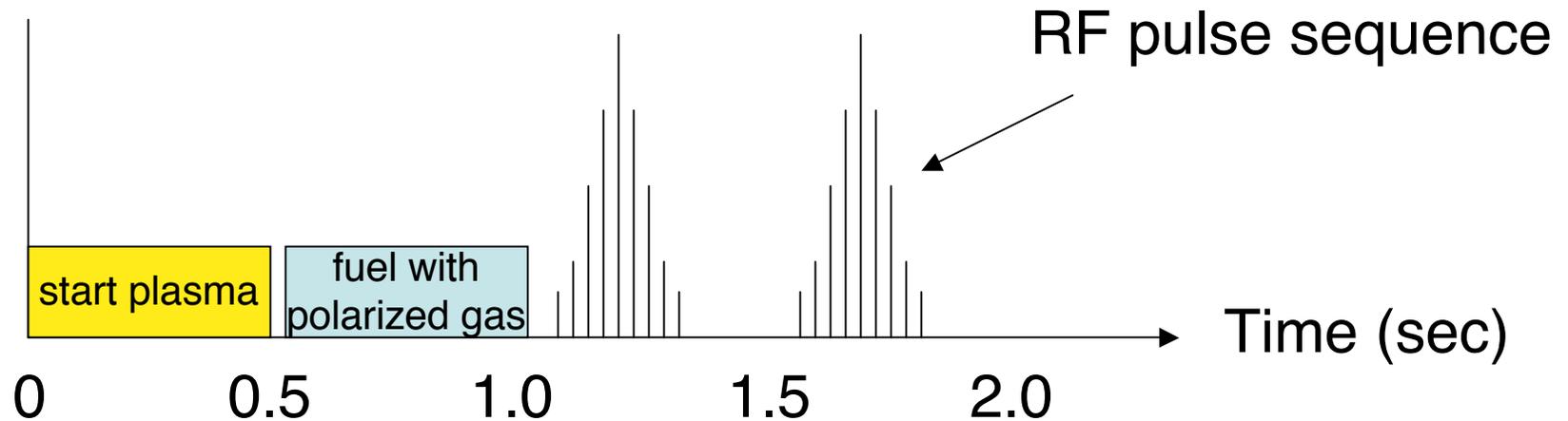
## Plasma NMR measurement:

- a) fuel plasma with highly polarized gas (e.g.  $^3\text{He}$ , H, or T)
- b) apply RF pulses to plasma using existing antennas
- c) measure NMR spectrum using same antennas
- d) vary B and RF pulse to scan resonance in plasma

## Wall Surface NMR measurement:

- a) apply RF pulses to empty vessel using existing antennas
- b) measure spectrum of spins using same antennas
- c) vary B and RF pulse to scan surface location

# Scenario for Tokamak NMR Measurement



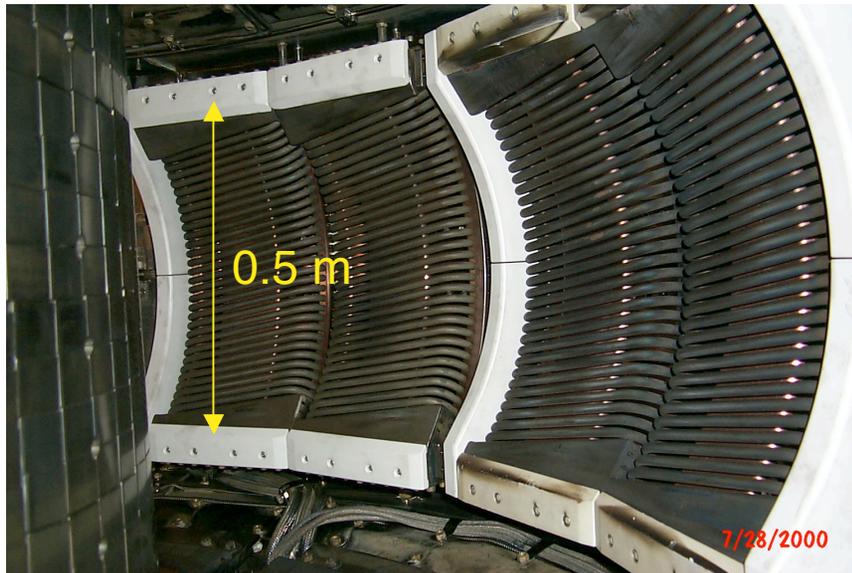
- Vary RF pulse amplitude to find maximum response
- Repeat to find timescale for plasma polarization decay
- Number of RF pulses limited by RF depolarization

# Use Existing RF Systems in Tokamaks

Note:  $\square_{\text{NMR}}/\square_{\text{ci}} = 2.8$  (H), 0.86 (D), 9 (T), 3.2 ( $^3\text{He}$ )

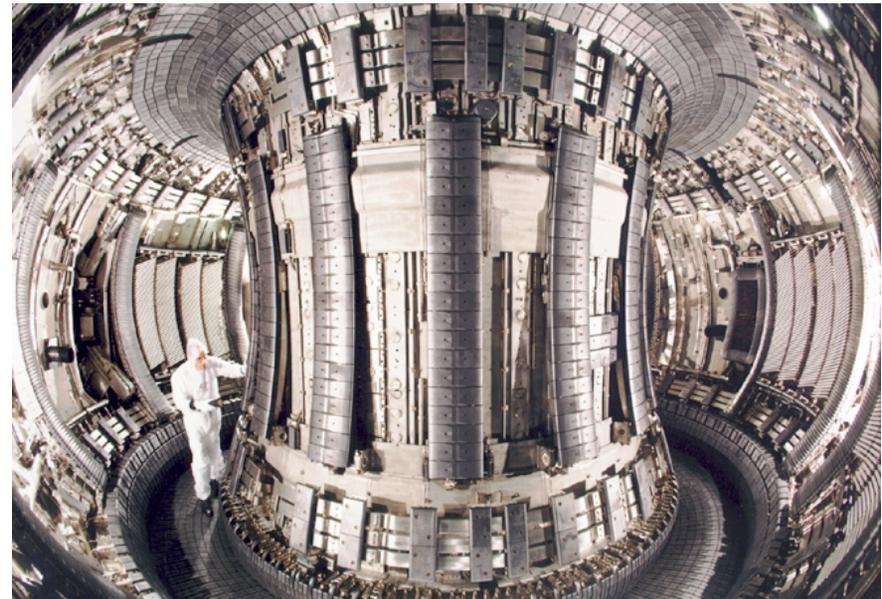
## Alcator C-Mod (MIT)

$B_{\text{toroidal}} \leq 8$  T,  $R=0.66$  m,  $a=0.23$  m  
 $f_{\text{ICRF}} \approx 40\text{-}80$  MHz  $P_{\text{ICRF}} \leq 5$  MW



## JET (Joint European Torus)

$B_{\text{toroidal}} \leq 3.8$  T,  $R=3.0$  m,  $a=1.0$  m  
 $f_{\text{ICRF}} \approx 25\text{-}100$  MHz  $P_{\text{ICRF}} \leq 20$  MW



# Potential Difficulties

## Plasma NMR measurement:

- Highly inhomogeneous toroidal magnetic field
- Short residence time of nuclei in a given volume
- Background due to atoms on walls
- RF propagation through plasmas
- RF noise from plasma

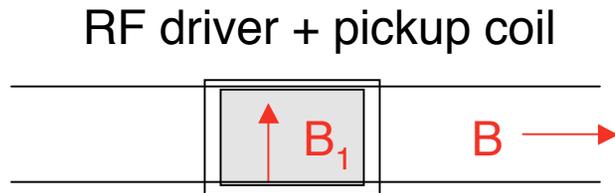
## Wall Surface NMR measurement:

- Small thickness of surface coatings ( $\approx 10 \mu\text{m}$ )
- RF skin effect in first-wall (e.g. carbon, steel)
- Coupling of RF radiation into empty vessel

# Estimates of Signal / Noise

- Plasma NMR
  - expected signal in uniform B
  - expected signal in tokamak B
  - expected NMR spectrum
  - background due to atoms on wall
  - estimate of RF noise from plasma
    - => expected S/N for plasma NMR
- Wall surface NMR
  - expected NMR signal from T on wall
  - expected NMR signal from lithium coating
  - overlap of NMR resonances on wall
    - => expected S/N for wall surface NMR

# Plasma NMR Signal in Uniform B



“Linear” Alcator C-Mod:

$$B=2.5 \text{ T (} f=80 \text{ MHz)}$$

$$n=2 \times 10^{20} \text{ m}^{-3}, T=1 \text{ keV}$$

➔ Assume fully polarized  $^3\text{He}$  plasma with spins rotated  $90^\circ$  to B by RF ( $B_1$ ) pulse

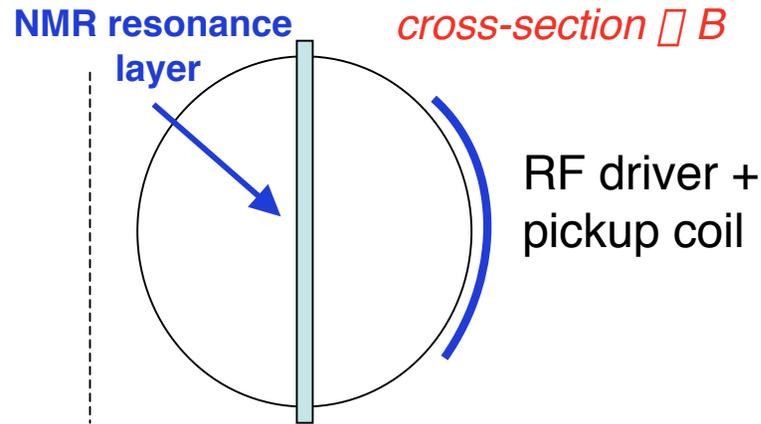
$$M_{\square} = n \mu_{3\text{He}} = n (2.1 \mu_{\text{N}}) \approx (2 \times 10^{20} \text{ m}^{-3}) (2.1 \cdot 5 \times 10^{-27} \text{ J/T}) = 2.1 \times 10^{-6} \text{ amp-turn/m}$$

$$B_{\square} = \mu_o M_{\square} = (4\pi \times 10^7 \text{ H/m})(2.1 \times 10^{-6} \text{ amp-turn/m}) = 2.6 \times 10^{-12} \text{ Tesla (assumes } \mu=\mu_o)$$

$$V_{\text{NMR}}(\text{volts}) = (N_{\text{coil}} A_{\text{coil}}) \omega_{\text{R}} B_{\square} = (0.3 \text{ m}^2)(2.5 \cdot 2 \times 10^8 \text{ rad/sec}) (2.6 \times 10^{-12} \text{ T}) \approx 4 \times 10^{-4} \text{ volts}$$

note: fractional polarization at room temperature  $f_{\text{pol}} \approx 2\mu_{3\text{He}} B / 2kT \approx 6 \times 10^{-6}$   
 fractional polarization at plasma temperature  $f_{\text{pol}} \approx (T_{\text{room}} / T_{\text{plasma}}) 6 \times 10^{-6} \approx 10^{-10}$

# Plasma NMR Signal in Tokamak B



## Actual Alcator C-Mod:

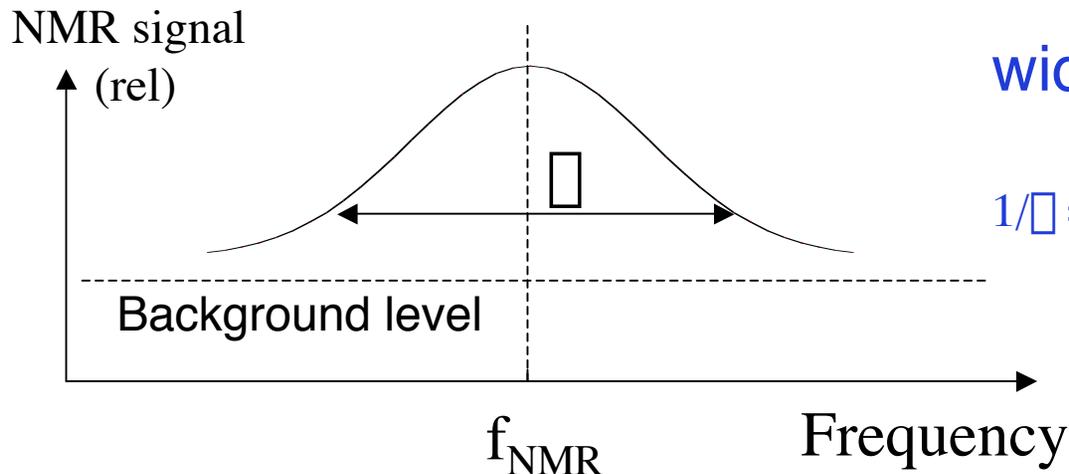
$$B = B_0(R_0/R) \text{ where } B_0 = 2.5 \text{ T}$$

$$f_{\text{NMR}} (^3\text{He}) = 80 \text{ MHz at } R_0 = 0.67 \text{ m}$$

$$n = 2 \times 10^{20} \text{ m}^{-3}, T = 1 \text{ keV}, \Delta_i = 0.1 \text{ cm}$$

- Width of NMR resonance layer:  $\Delta R/R \approx \Delta B/B \approx \Delta f_{\text{driver}}/f_{\text{NMR}}$
- Driver  $\Delta f_{\text{driver}}$  depends on RF pulse width:  $\Delta f_{\text{driver}} \Delta_{\text{RF}} \approx 1$
- For  $\Delta R \approx 1 \text{ cm}$  ( $\approx 10 \Delta_i$ ) and  $f_{\text{NMR}} \approx 80 \text{ MHz} \Rightarrow \Delta_{\text{RF}} \approx 1 \mu\text{sec}$
- For “filling factor”  $\approx 2\%$  of uniform B case:  $V_{\text{NMR}} \approx 10^{-5} \text{ volts}$

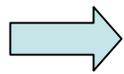
# Expected NMR Spectrum of a Tokamak



width of NMR spectrum

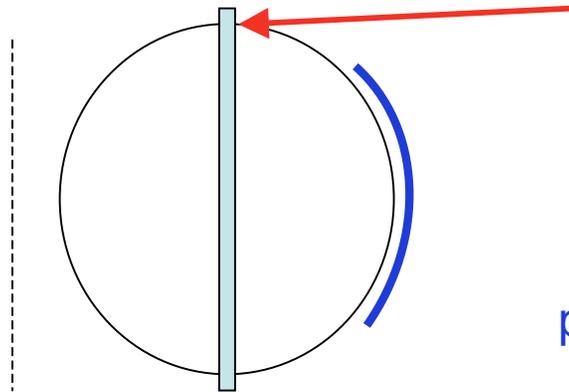
$1/\Delta = \Delta$  all dephasing processes

$\Delta_B$	magnetic field gradients	$(\Delta B/B) \Delta_{\text{NMR}}$	$\approx 1 \text{ MHz}$
$\Delta_{\text{pol}}$	motion along poloidal field	$(v_i/\Omega R)(B_{\text{pol}}/B)$	$\approx 1 \text{ MHz}$
$\Delta_{\Omega}$	ion cyclotron motion	$(\Omega_i/R) \Delta_{\text{NMR}}$	$\approx 0.3 \text{ MHz}$
$\Delta_r$	diffusion across B	$D/\Omega R^2$	$\approx 10^4 \text{ Hz}$
$\Delta_v$	thermal Doppler broadening	$(v_i/c) \Delta_{\text{NMR}}$	$\approx 10^3 \text{ Hz}$
$\Delta_{\delta B}$	magnetic field fluctuations	$(\Delta B_{\delta}/B)^2 \Delta_{\text{NMR}}$	$\approx 10 \text{ Hz}$



Spectrum dominated by  $\Delta_B$  and poloidal motion  
*Spectral width much broader than solids ( $\approx 1 \text{ Hz}$ )*

# Background Due to Fuel Atoms on Walls



NMR resonance interacts with nuclei stuck on surface of wall

1 monolayer on wall surface  $\approx 10^{15} - 10^{16}$  atoms/cm<sup>2</sup>

plasma "na"  $\approx (2 \times 10^{14} \text{ cm}^{-3})(50 \text{ cm}) \approx 10^{16}$  atoms/cm<sup>2</sup>

<sup>3</sup>He probably has  $\ll 1$  monolayer  $\Rightarrow$  negligible

However H, D, T up to  $\approx 10^4$  monolayers  $\approx 10^{19} - 10^{20}$  atoms/cm<sup>2</sup>

$\Rightarrow$  But if only  $f_{\text{pol}} \approx 10^{-5}$  wall atoms polarized room temperature, then wall background still  $<$  plasma signal

# RF “ICE” Emission from Plasma

Spectra of RF emission as seen by a JET antenna (Cottrell & Dendy, PRL ‘88)

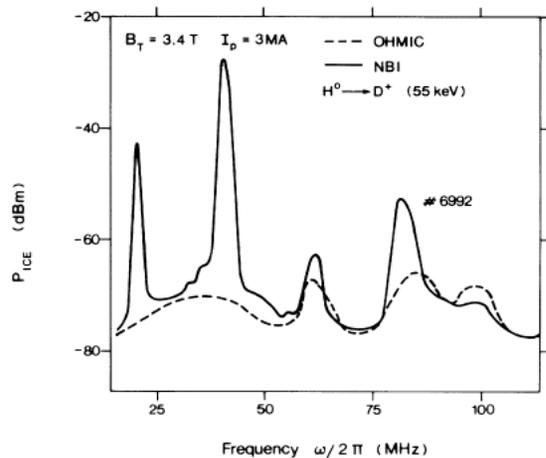


FIG. 1. ICE spectrum measured before and during 4-MW  $H^0$  neutral-beam co-injection into a  $D^+$  limiter plasma.

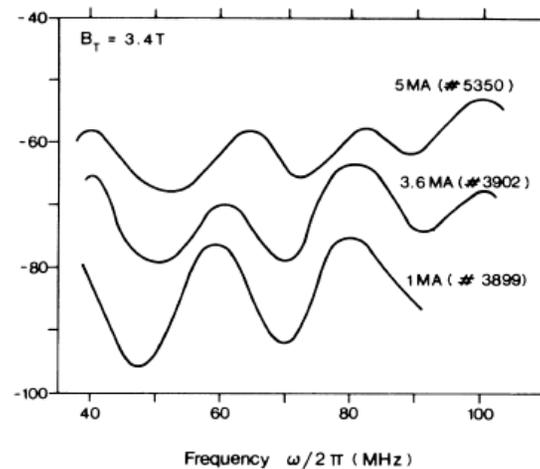


FIG. 2. ICE spectrum from three Ohmically heated deuterium limiter discharges with constant toroidal field but with different plasma currents.

“ICE” lines at  $n \approx n_{ci}$  of beam ions or fusion products [McClements et al, PRL ‘99]

For low power Ohmic plasmas in JET the measured ICE was as low as  $\approx -80$  dBm

$$\square P_{RF} \approx 10^{-11} \text{ Watt into } \approx 1 \text{ m}^2 \text{ antenna}$$



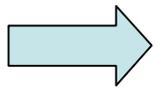
Arunasalam tried to explain ICE as nuclear spin-flip emission ! (Nucl. Fusion ‘94)

## Estimated S/N for Plasma NMR

- Plasma NMR signal  $V_{\text{NMR}} \approx 10^{-5}$  volts (Alcator C-Mod  $^3\text{He}$ )
  - Assume plasma noise  $P_{\text{noise}} \approx -80$  dBm  $\approx 10^{-11}$  W (approx.)
  - For  $R = 50 \text{ } \Omega$  antenna:  $V_{\text{noise}} \approx (P_{\text{noise}} R)^{1/2} \approx 2 \times 10^{-5}$  volts  
 $\Rightarrow$  Plasma S/N  $\approx 0.5$
  - Detector noise ( $\Delta f = 1$  MHz):  $V_{\text{det}} = (4k_{\text{B}} T_{\text{room}} R \Delta f)^{1/2} \approx 10^{-6}$  volts  
 $\Rightarrow$  Detector S/N  $\approx 10$
- $\Rightarrow$  S/N could be increased by  $\approx \times 4$  with multiple pulses/shot

# NMR Wave Propagation in Tokamaks

- RF heating antennas excite the fast magnetosonic wave and a small component of the fast wave field is properly polarized to couple to (negative)  $^3\text{He}$  nuclear spin
- NMR frequency of  $^3\text{He}$  is  $\approx 3\omega_{^3\text{He}}$  and  $\approx 4\omega_{\text{D}}$ , so fast wave damping at harmonics is a possible loss mechanism



RF waves from antenna can couple to nuclear spin  
Spins can couple to fast wave and back to antenna

BUT, expected  $\Delta B_{\parallel}$  from NMR  $\approx 10^{-8}$  Gauss, while expected plasma-driven magnetic fluctuations  $\approx 1$  Gauss !?

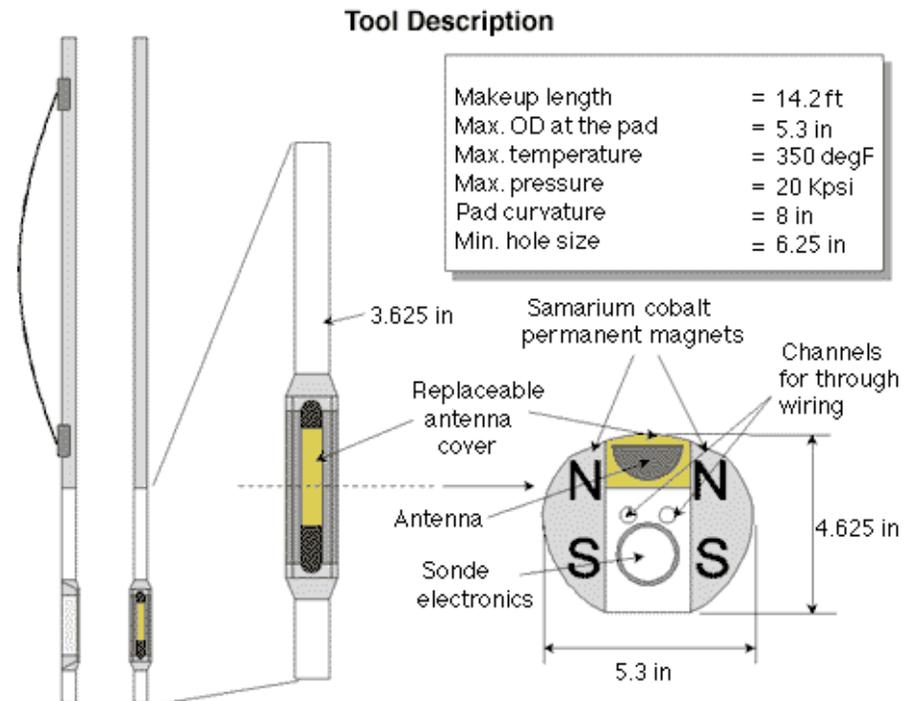
*=> why can't we see the internal plasma  $\Delta B_{\parallel}$  without NMR ?*

# Wall Surface NMR

- Measure absolute wall surface composition without plasma
- Maybe can use TF (VF?) field + RF antenna (near field only)
- Maybe can insert NMR probe into vessel with remote arm
- Should work up to RF skin depth,  $\approx 10^{-2}$  cm at 100 MHz

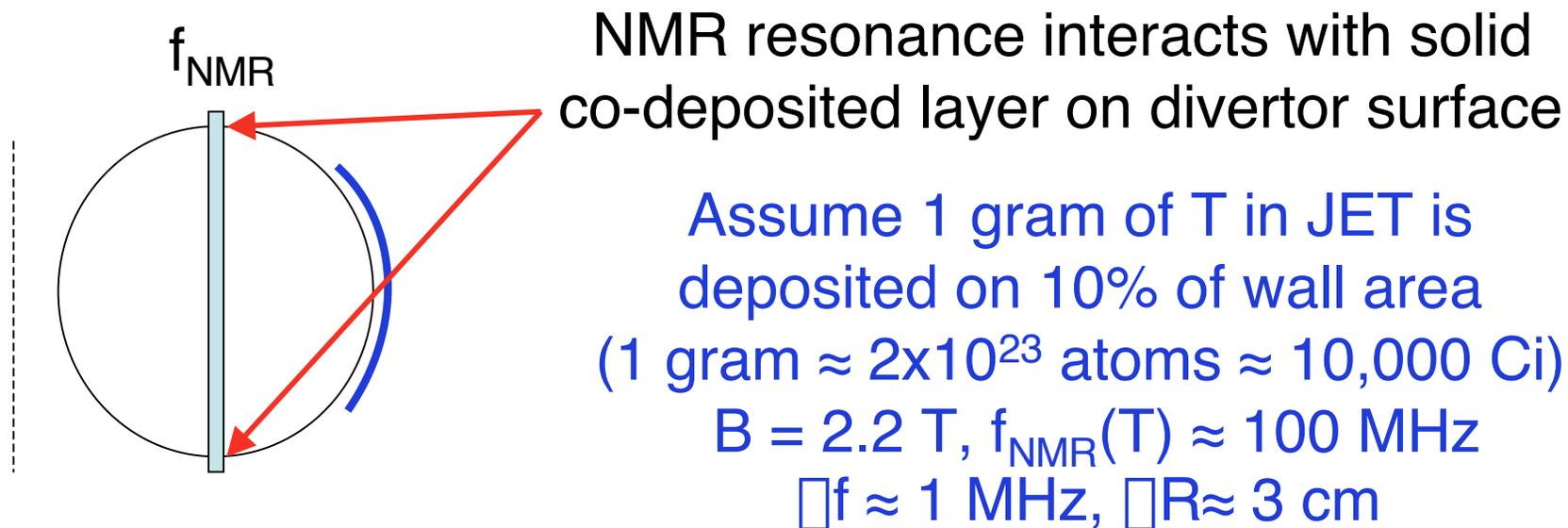
Tritium detection threshold in lab NRM  $\approx 0.1$  mCi,  $\approx 10^{-8}$  of tritium levels inside JET

“Inside out” NMR used for oil exploration (Schlumberger)



<http://www.slb.com/seed/en/watch/nmr/tool.htm>

## Measurement of Tritium on Surface



$$N_{\text{T}} \approx (2 \times 10^{23} \text{ atoms}) / (2 \times 10^5 \text{ cm}^2) \approx 10^{18} \text{ atoms/cm}^2 \text{ on surface}$$

$$\text{Assume } f_{\text{pol}} = (\mu B_{\text{T}} / \mu B_{3\text{He}}) (6 \times 10^{-6}) \approx 7 \times 10^{-6} \text{ polarized}$$

Assume this is equivalent to distributed density in plasma of

$$n \approx (10^{18} \text{ atoms/cm}^2) (7 \times 10^{-6}) / 300 \text{ cm} \approx 2 \times 10^{10} \text{ atoms/cm}^3$$

## Estimated S/N for Tritium on Surface

- Assume detection of 3 cm radius with JET antenna area 1 m<sup>2</sup>
- Signal  $V_{\text{NMR}} \propto (NA_{\text{coil}})n\mu\hbar B$  (scaling from <sup>3</sup>He case in Alcator)

$$V_{\text{NMR}} \propto (1.0 \text{ m}^2/0.3 \text{ m}^2)(2 \times 10^{10} \text{ cm}^{-3}/2 \times 10^{14} \text{ cm}^{-3})(1.4)(100 \text{ MHz}/80 \text{ MHz})(2.2 \text{ T}/2.5 \text{ T})$$

$$V_{\text{NMR}} \approx (5 \times 10^{-5}) (10^{-5} \text{ volts}) \approx 5 \times 10^{-9} \text{ volts}$$

- Detector noise ( $\Delta f = 1 \text{ MHz}$ ):  $V_{\text{det}} = (4k_{\text{B}}T_{\text{room}}R\Delta f)^{1/2} \approx 10^{-6} \text{ volts}$

$$\Rightarrow S/N \approx 5 \times 10^{-3}$$

This could be improved by multiple RF pulsing during a 10 sec long JET shot (depending on spin relaxation time in layer)

# Measurement of Lithium on Surface

$f_{\text{NMR}} = 63 \text{ MHz}$  for  ${}^7\text{Li}$  NMR at maximum  $B=3.8 \text{ T}$  in JET

$N_{\text{Li}} \approx 5\%(10^4 \text{ monolayers})(5 \times 10^{15} \text{ atoms/cm}^2) \approx 2 \times 10^{18} \text{ atoms/cm}^2$

$f_{\text{pol}} = (\mu_{\text{Li}}/\mu_{\text{He}})(6 \times 10^{-6}) \approx (3.2/2.1)(3.8/2.2) 6 \times 10^{-6} \approx 1.5 \times 10^{-5} \text{ polarized}$

Assume this is equivalent to distributed density in plasma of:

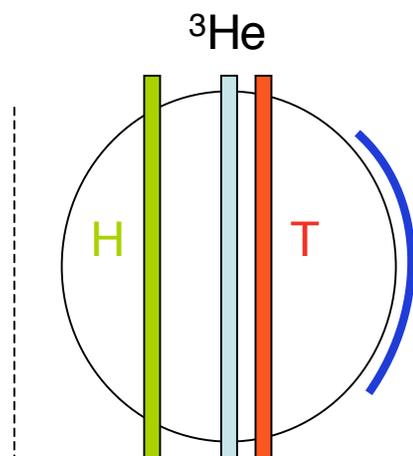
$n \approx (2 \times 10^{18} \text{ atoms/cm}^2)(1.5 \times 10^{-5})/300 \text{ cm} \approx 3 \times 10^{11} \text{ atoms/cm}^3$

Signal  $V_{\text{NMR}} \propto (NA_{\text{coil}})n\mu \square B$  (scaling from T case in Alcator)

$V_{\text{NMR}}(\text{Li}) \approx (150)(3.2/3.0)(60/100)(3.8/2.2)V_{\text{NMR}}(\text{T}) \approx 2 \times 10^{-8} \text{ volts}$

$\Rightarrow S/N \approx 150(5 \times 10^{-4}) \approx 7 \times 10^{-2}$

# Overlap of NMR Resonances



If antenna set at  $f=80$  MHz at 2.5 T:

<u>Species</u>	<u>Resonance layer</u>
$^3\text{He}$	$R_0$
H	$0.76 R_0$
D	$6 R_0$
T	$1.06 R_0$
$^6\text{Li}$	$5 R_0$
$^7\text{Li}$	$2 R_0$
$^{13}\text{C}$	$4 R_0$

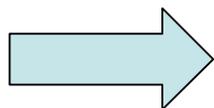
Problem:

- H and T will probably overlap - indistinguishable ?
- D and  $^6\text{Li}$  will probably overlap - indistinguishable ?

*=> may be able to avoid overlap with localized NMR probe*

# NMR of Surface Tile Removed from Vessel

- First wall tiles are routinely removed from the tokamak for surface analysis (which can still be difficult)
- Could wall samples be analyzed in an NMR spectrometer ?
- For DT in JET or TFTR, tile has typically 5 mCi/cm<sup>2</sup> of tritium on surface (within upper 10<sup>4</sup> -10<sup>5</sup> monolayers)
- NMR is routinely done with 0.1-10 mCi/sample (Evans et al, "Handbook of Tritium NMR Spectroscopy", 1985, p. xii)



Looks like NMR could be done OK outside vessel

## Summary of S/N Estimates

- For  $^3\text{He}$  plasma NMR in Alcator C-Mod  
S/N  $\approx 0.5 - 10$  for one NMR pulse  
S/N  $\approx 2 - 40$  for  $\approx 20$  NMR pulses/shot
- For tritium or lithium on wall surface in JET  
S/N  $\approx 10^{-2}$  for one NMR pulse  
S/N  $\approx 1$  for  $10^4$  NMR pulses during long TF
- For “outside the tokamak” wall tile surface analysis  
S/N  $\gg 1$  inside NMR spectrometer

## Conclusions

- Measuring NMR response of fully polarized  $^3\text{He}$  plasma in Alcator C-Mod seems **marginally feasible**
- Measurement of internal magnetic fluctuations via decay of polarization is **difficult** due to uncertainties about propagation of RF through plasma and ICE emission
- Measurement of low-Z composition of wall (without plasma) seems **marginally feasible** but **potentially useful**, e.g. for assessment of tritium inventory in JET or ITER

## Possible Directions

- Analyze TFTR tiles for H content using NMR
  - use small sample in NMR spectrometer
  - compare with other surface analysis methods
  - develop *in situ* wall surface T diagnostic for ITER
- Do initial plasma experiments on a linear machine
  - much lower  $\square B$  => much larger NMR S/N ratio
  - less RF noise and simpler RF propagation physics
  - try NMR “tagging”, flow, and diffusion experiments
- Try spin-polarized fueling on a toroidal machine
  - try to measure NMR of polarized  $^3\text{He}$  in Alcator
  - if successful, try spin-polarized T fueling in JET
  - pursue x1.5 increase in DT fusion reaction rate